

# Spectral properties of random banded Hessenberg matrices

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## Banded lower Hessenberg matrices

In this talk we consider **banded lower Hessenberg matrices**

$$H_n = \begin{pmatrix} a_1^{(0)} & 1 & & & & 0 \\ a_1^{(1)} & a_2^{(0)} & 1 & & & \\ \vdots & \ddots & \ddots & \ddots & & \\ a_1^{(p)} & & \ddots & \ddots & \ddots & \\ & \ddots & & \ddots & \ddots & 1 \\ 0 & & a_{n-p}^{(p)} & \cdots & a_{n-1}^{(1)} & a_n^{(0)} \end{pmatrix}$$

with 1's in the first superdiagonal.

The number of subdiagonals is  $p \geq 1$ , arbitrary fixed positive integer.

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$a_j^{(k)}$  is located in the  $k$ th subdiagonal ( $k = 0$  for the main diagonal) and the  $j$ th column.

The characteristic polynomials  $Q_n(z) = \det(zI_n - H_n)$  satisfy the  **$(p + 2)$ -term recurrence relation**

$$zQ_n(z) = Q_{n+1}(z) + a_{n+1}^{(0)}Q_n(z) + a_n^{(1)}Q_{n-1}(z) + \cdots + a_{n-p+1}^{(p)}Q_{n-p}(z).$$

The roots of  $Q_n$  (eigenvalues of  $H_n$ ) are in general spread in the complex plane in a rather arbitrary way.

## Some notable examples

Examples of such polynomials are **multiple orthogonal polynomials** with respect to **Angelesco systems** or **Nikishin systems** of  $p$  measures on the real line, or on a symmetric star with center at the origin in the complex plane with rotationally symmetric generating measures.

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In the case of Angelesco systems on the real line: Take  $\mu_1, \dots, \mu_p$  positive measures with infinite support  $\text{supp}(\mu_k) \subset [a_k, b_k]$ , such that  $[a_k, b_k] \cap [a_j, b_j] = \emptyset$  for  $k \neq j$ . Then  $Q_n$  is the  $n$ th-degree monic polynomial satisfying

$$\int Q_n(x) x^k d\mu_j(x) = 0, \quad k = 0, \dots, \left\lfloor \frac{n-j}{p} \right\rfloor, \quad 1 \leq j \leq p.$$

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Then the sequence  $(Q_n(z))_{n=0}^{\infty}$  satisfies

$$zQ_n(z) = Q_{n+1}(z) + a_{n+1}^{(0)}Q_n(z) + a_n^{(1)}Q_{n-1}(z) + \dots + a_{n-p+1}^{(p)}Q_{n-p}(z), \quad n \geq 0,$$

with initial conditions  $Q_0 \equiv 1, Q_{-1} \equiv \dots \equiv Q_{-p} \equiv 0$ .

## Banded Hessenberg operators and resolvent functions

Let  $\mathbf{a}^{(k)} = (a_n^{(k)})_{n=1}^{\infty}$ ,  $0 \leq k \leq p$ , be bounded sequences of complex numbers, and construct the bounded operator  $H$  on  $\ell^2(\mathbb{N})$  with matrix representation

$$H = \begin{pmatrix} a_1^{(0)} & 1 & & & & & 0 \\ \vdots & \ddots & \ddots & & & & \\ a_1^{(p)} & \ddots & \ddots & \ddots & & & \\ & \ddots & \ddots & \ddots & & & \\ & & a_{n-p}^{(p)} & \ddots & a_n^{(0)} & \ddots & \\ 0 & & & \ddots & \ddots & \ddots & \end{pmatrix}.$$

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Let  $\{\mathbf{e}_j\}_{j=1}^{\infty}$  be the standard basis in  $\ell^2(\mathbb{N})$ , with inner product  $\langle \cdot, \cdot \rangle$ . Let

$$\phi_j(z) = \langle (z - H)^{-1} \mathbf{e}_j, \mathbf{e}_1 \rangle = \sum_{n=0}^{\infty} \frac{\langle H^n \mathbf{e}_j, \mathbf{e}_1 \rangle}{z^{n+1}}, \quad 1 \leq j \leq p.$$

We also define  $\phi_0 \equiv 1$ . We have defined a map

$$\mathcal{A} = (\mathbf{a}^{(0)}, \mathbf{a}^{(1)}, \dots, \mathbf{a}^{(p)}) \mapsto \Phi = (\phi_0, \phi_1, \dots, \phi_p)$$

which we indicate by writing  $\Phi = \Phi(\mathcal{A})$ .



## Random matrices

Assumptions:

- 1) Let  $\mu_k$ ,  $0 \leq k \leq p$ , be a collection of  $p + 1$  Borel probability measures with compact support in  $\mathbb{C}$ .

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Construct the infinite banded Hessenberg matrix

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Let  $H_n$  be the  $n \times n$  principal truncation of  $H$ :

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Let  $\{\lambda_{i,n}\}_{i=1}^n$  denote the eigenvalues of  $H_n$ , counting multiplicities, and let

$$\sigma_n := \frac{1}{n} \sum_{i=1}^n \delta_{\lambda_{i,n}}. \quad (\text{ESD of } H_n)$$

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$\mathbb{E}\sigma_n$  is the deterministic probability measure defined via duality by

$$\int f d\mathbb{E}\sigma_n = \mathbb{E} \left( \int f d\sigma_n \right).$$

Main questions of interest:

- 1) Is the sequence of average measures  $\mathbb{E}\sigma_n$  weakly convergent (in the weak-star topology)?
- 2) If so, what is the limit, and how is it related to the distributions  $\{\mu_k\}_{k=0}^p$ ?

We have a partial answer to these questions, proving the existence of the limits of the average moments

$$\lim_{n \rightarrow \infty} \int z^\ell d\mathbb{E}\sigma_n(z) = \lim_{n \rightarrow \infty} \mathbb{E} \left( \frac{1}{n} \sum_{i=1}^n \lambda_{i,n}^\ell \right) =: \omega_\ell, \quad \ell \geq 0, \quad (1)$$

and we can describe the **moment generating function**

$$\sum_{\ell=0}^{\infty} \frac{\omega_\ell}{z^{\ell+1}}$$

in terms of resolvent functions of the operator  $H$  and an extension of this operator to  $\ell^2(\mathbb{Z})$ .



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In certain models of bi-diagonal random Hessenberg matrices with eigenvalues on a symmetric starlike set, we can also prove weak convergence of  $\mathbb{E}\sigma_n$  to a probability measure on the set.

Recall the map

$$\mathcal{A} = (\mathbf{a}^{(0)}, \mathbf{a}^{(1)}, \dots, \mathbf{a}^{(p)}) \mapsto \Phi = (\phi_0, \phi_1, \dots, \phi_p)$$

which we indicate by writing  $\Phi = \Phi(\mathcal{A})$ , where  $\phi_0 \equiv 1$  and

$$\phi_j(z) := \langle (z - H_{\mathcal{A}})^{-1} \mathbf{e}_j, \mathbf{e}_1 \rangle, \quad 1 \leq j \leq p.$$

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### Theorem

*Let  $\mathcal{A} = (\mathbf{a}^{(0)}, \dots, \mathbf{a}^{(p)})$  and  $\mathcal{B} = (\mathbf{b}^{(0)}, \dots, \mathbf{b}^{(p)})$  be two independent collections of random sequences with corresponding distributions  $(\mu_0, \dots, \mu_p)$ . Let  $\Phi(\mathcal{A}) = (\phi_0, \dots, \phi_p)$  and  $\Psi(\mathcal{B}) = (\psi_0, \dots, \psi_p)$ .*

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$$\lim_{n \rightarrow \infty} \int z^\ell d\mathbb{E}\sigma_n(z) = \mathbb{E}([W]_{\ell+1}),$$

where  $[W]_{\ell+1}$  is the coefficient of  $z^{-\ell-1}$  in the Laurent series expansion at  $\infty$  of

$$W(z) = \frac{1}{z - \sum_{k=0}^p \sum_{j=0}^k \alpha_j^{(k)} \phi_{k-j}(z) \psi_j(z)}.$$

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The expression  $\sum_{k=0}^p \sum_{j=0}^k \alpha_j^{(k)} \phi_{k-j}(z) \psi_j(z)$  is the triangular bilinear form:

$$(\psi_0 \ \psi_1 \ \psi_2 \ \dots \ \psi_p) \begin{pmatrix} \alpha_0^{(0)} & \alpha_0^{(1)} & \alpha_0^{(2)} & \dots & \alpha_0^{(p)} \\ \alpha_1^{(1)} & \alpha_1^{(2)} & \dots & \alpha_1^{(p)} & \\ \alpha_2^{(2)} & \dots & \alpha_2^{(p)} & & \\ \vdots & \ddots & & & \\ \alpha_p^{(p)} & & & & 0 \end{pmatrix} \begin{pmatrix} \phi_0 \\ \phi_1 \\ \phi_2 \\ \vdots \\ \phi_p \end{pmatrix}.$$

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The function  $W(z)$ , analytic in a nbhd of  $\infty$ , is a resolvent function of a two-sided operator on  $\ell^2(\mathbb{Z})$ .

The operator in  $\ell^2(\mathbb{Z})$  is obtained extending the sequences in  $\mathcal{A} = (a^{(0)}, a^{(1)}, \dots, a^{(p)})$  from  $\mathbb{N}$  to  $\mathbb{Z}$  using the sequences in  $\mathcal{B} = (b^{(0)}, \dots, b^{(p)})$  and the array  $\alpha = (\alpha_j^{(k)})_{0 \leq j \leq k \leq p}$ : For  $n \leq 0$ , we define

$$a_n^{(k)} = \begin{cases} \alpha_{-n}^{(k)} & \text{if } -k \leq n \leq 0, \\ b_{-n-k}^{(k)} & \text{if } n \leq -k - 1. \end{cases}$$





## Some key ideas in the proof

1) **Hermite-Padé property:** Let

$$Q_n(z) = \det(zI_n - H_n),$$
$$Q_{n,j}(z) = \det((zI_n - H_n)^{[j]}), \quad 1 \leq j \leq p, \quad n \geq 0,$$

where  $(zI_n - H_n)^{[j]}$  is the submatrix of  $zI_n - H_n$  obtained deleting the first  $j$  rows and columns.

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Kalyagin (1995) proved that

$$\left( \frac{Q_{n,1}}{Q_n}, \frac{Q_{n,2}}{Q_n}, \dots, \frac{Q_{n,p}}{Q_n} \right)$$

is a Hermite-Padé approximant at infinity for the system of resolvent functions  $(\phi_1, \phi_2, \dots, \phi_p)$  associated with the operator  $H$ :

$$Q_n(z)\phi_j(z) - Q_{n,j}(z) = O\left(\frac{1}{z^{n_j+2}}\right), \quad z \rightarrow \infty, \quad 1 \leq j \leq p,$$

where  $n_j = \lfloor (n-j)/p \rfloor$ .

2) Let  $M = (m_{i,j})_{i,j \in \mathbb{Z}}$  be the bi-infinite banded Hessenberg matrix defined before. For  $n \geq 0$ , let

$$M_{2n+1} = (m_{i,j})_{-n \leq i,j \leq n}.$$

Then for every  $-n \leq j \leq n$ ,

$$(zI_{2n+1} - M_{2n+1})^{-1}(j,j) - w_j(z) = O\left(\frac{1}{z^{n-|j|+3+\lfloor(n-|j|)/p\rfloor}}\right), \quad z \rightarrow \infty \quad (2)$$

where

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(2) follows from the Hermite-Padé property and the identity

$$(zI_{2n+1} - M_{2n+1})^{-1}(j,j) = \frac{q_{n-j}^+(z)q_{n+j}^-(z)}{q_{2n+1}(z)}$$

where

$$\begin{aligned} q_{2n+1}(z) &= \det(zI_{2n+1} - M_{2n+1}) \\ q_\ell^+(z) &= \det((zI_{2n+1} - M_{2n+1})^{[2n+1-\ell]}) \\ q_\ell^-(z) &= \det((zI_{2n+1} - M_{2n+1})_{[2n+1-\ell]}) \end{aligned}$$

In fact, the key relations are

$$\begin{aligned}
 (zI_{2n+1} - M_{2n+1})^{-1}(j, j) &= \frac{q_{n-j}^+(z)q_{n+j}^-(z)}{q_{2n+1}(z)} \\
 &= \frac{1}{z - a_j^{(0)} - \sum_{t=1}^p \sum_{k=0}^t a_{j-k}^{(t)} \frac{q_{n-j+k-t}^+(z)}{q_{n-j}^+(z)} \frac{q_{n+j-k}^-(z)}{q_{n+j}^-(z)}} \\
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where  $\phi_{j,t-k}^+(z)$  and  $\phi_{j,k}^-(z)$  are the resolvent functions of certain restrictions of the operator  $M$  on  $\ell^2(\mathbb{N})$ .

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3) The connection with the eigenvalues  $\lambda_1, \lambda_2, \dots, \lambda_{2n+1}$  of  $M_{2n+1}$  is

$$\sum_{j=-n}^n (zI_{2n+1} - M_{2n+1})^{-1}(j, j) = \sum_{k=1}^{2n+1} \frac{1}{z - \lambda_k}$$

From the independence between the elements  $\alpha_j^{(k)}$ ,  $\phi_{k-j}(z)$ , and  $\psi_j(z)$  in the formula

$$W(z) = \frac{1}{z - \sum_{k=0}^p \sum_{j=0}^k \alpha_j^{(k)} \phi_{k-j}(z) \psi_j(z)}$$

one can easily obtain a relation (in the form of a series) between  $\mathbb{E}(W(z))$ , moments of the distributions  $\mu_0, \mu_1, \dots, \mu_p$ , and joint moments of the random vector  $(\phi_1(z), \dots, \phi_p(z))$ :

$$g_{(n_1, \dots, n_p)}(z) = \mathbb{E} \left( \prod_{k=1}^p \phi_k(z)^{n_k} \right).$$



## The distribution of $(\phi_1(z), \dots, \phi_p(z))$

For  $z$  large enough, the random vector  $(\phi_1(z), \dots, \phi_p(z))$  is well-defined (resolvent functions of the infinite matrix  $H$  on  $\ell^2(\mathbb{N})$ ). Let  $\sigma_z$  be the distribution of this vector (probability measure on  $\mathbb{C}^p$ ). We don't know what it is, but it satisfies an **invariance principle**.

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Let  $H_1$  be the infinite one-sided matrix obtained by removing the first row and the first column of the matrix  $H$ , and let

$$\begin{aligned}\phi_j(z) &= \langle (z - H)^{-1} \mathbf{e}_j, \mathbf{e}_1 \rangle & 1 \leq j \leq p \\ \phi_{1,j}(z) &= \langle (z - H_1)^{-1} \mathbf{e}_j, \mathbf{e}_1 \rangle & 1 \leq j \leq p\end{aligned}$$

then for  $z$  large enough

$$\phi_1(z) = \frac{1}{z - a_1^{(0)} - \sum_{k=1}^p a_1^{(k)} \phi_{1,k}(z)}, \quad (3)$$

$$\phi_j(z) = \frac{\phi_{1,j-1}(z)}{z - a_1^{(0)} - \sum_{k=1}^p a_1^{(k)} \phi_{1,k}(z)}, \quad 2 \leq j \leq p, \quad (4)$$

which when iterated gives a **vector continued fraction** expansion of  $(\phi_1(z), \dots, \phi_p(z))$  (vector analogue of the Jacobi continued fraction for OPs).

By the i.i.d. assumptions, the vectors  $\Phi(z) := (\phi_1(z), \phi_2(z), \dots, \phi_p(z))$  and  $\Phi_1(z) := (\phi_{1,1}(z), \phi_{1,2}(z), \dots, \phi_{1,p}(z))$  clearly have the same distribution  $\sigma_z$ , and the vectors  $(a_1^{(0)}, \dots, a_1^{(\rho)})$  and  $\Phi_1(z)$  are independent.

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We will write  $\mathbf{t} = (t_0, \dots, t_p)$ ,  $\mathbf{x} = (x_1, \dots, x_p)$ . Let  $\lambda_z : \mathcal{D}_z \rightarrow \mathbb{C}^p$  be the function

$$\lambda_z(\mathbf{t}, \mathbf{x}) := \left( \frac{1}{z - t_0 - \sum_{k=1}^p t_k x_k}, \frac{x_1}{z - t_0 - \sum_{k=1}^p t_k x_k}, \dots, \frac{x_{p-1}}{z - t_0 - \sum_{k=1}^p t_k x_k} \right).$$

Then the relations (3)–(4) mean  $\Phi(z) = \lambda_z((a_1^{(0)}, \dots, a_1^{(p)}), \Phi_1(z))$ , and the distribution of  $((a_1^{(0)}, \dots, a_1^{(p)}), \Phi_1(z))$  is  $\mu \times \sigma_z$ , so

$$\sigma_z = (\lambda_z)_*(\mu \times \sigma_z)$$

(push-forward of  $\mu \times \sigma_z$  under  $\lambda_z$ ).

This means that for any  $f : \mathbb{C}^p \rightarrow \mathbb{C}$  in  $L^1(\sigma_z)$ , we have

$$\int f d\sigma_z = \iint f \circ \lambda_z d(\mu \times \sigma_z).$$

## Literature on random matrices

Most of the literature on random matrices is devoted to the study of ensembles of Hermitian/real-symmetric or unitary/orthogonal matrices. The great advantage of such ensembles is that the eigenvalues are located on the real line or the unit circle. Important examples are GUE, GOE, GSE, CUE.

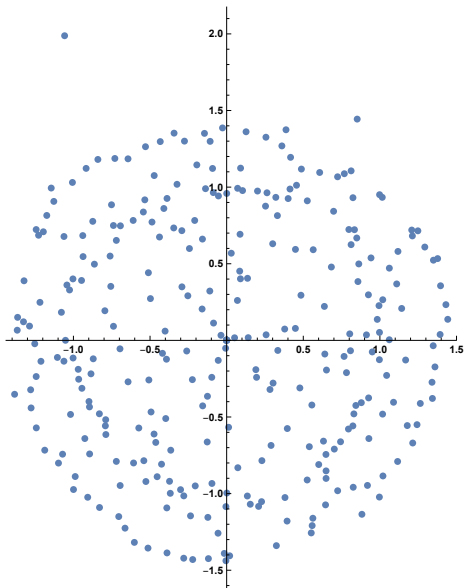
In the class of random banded matrices, most works in the literature also assume symmetry (see e.g. works of Bourgade, Fyodorov, M. Shcherbina, T. Shcherbina, Sodin, Spencer, and others).

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Random banded Hessenberg matrices are on the contrary non-symmetric, and so it is difficult in general to locate the eigenvalues, so the spectral analysis presents new challenges. We have found that Hermite-Padé approximation can be an effective tool for this analysis.



**Figure:** Eigenvalues of  $H_n$ ,  $n = 300$ ,  $p = 2$ ,  $a_n^{(0)} = a_n^{(1)} = 0$  for all  $n$ ,  $(a_n^{(2)})_{n=1}^{\infty}$  is i.i.d. with uniform distribution on the unit circle.



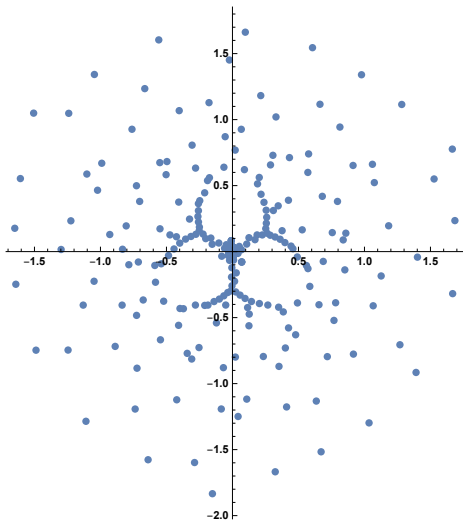


Figure: Eigenvalues of  $H_n$ ,  $n = 300$ ,  $p = 2$ ,  $a_n^{(0)} = a_n^{(1)} = 0$  for all  $n$ ,  $(a_n^{(2)})_{n=1}^{\infty}$  is i.i.d. with uniform distribution on the arc of the unit circle  $\{e^{i\theta} : \frac{\pi}{3} \leq \theta \leq \frac{2\pi}{3}\}$ .

## References

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