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# Smooth TiO<sub>2</sub> Thin Films Grown by Aqueous Spray Deposition for Long-Wave Infrared Applications

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ABSTRACT

Self-assembled TiO<sub>2</sub> films deposited by aqueous-spray deposition were investigated to evaluate morphology, crystalline phase, and infrared optical constants. The Anatase nanocrystalline film had ~10 nm characteristic surface roughness sparsely punctuated by defects of not more than 200 nm amplitude. The film is highly transparent throughout the visible to wavelengths of 12 µm. The indirect band gap was determined to be 3.2 eV. Important for long-wave infrared applications is that dispersion in this region is weak compared with the more commonly used dielectric SiO<sub>2</sub> for planar structures. An example application to a metal-insulator-metal resonant absorber is presented. The low-cost, large-area, atmosphericpressure, chemical spray deposition method allows conformal fabrication on flexible substrates for long-wave infrared photonics.

### **INTRODUCTION**

Titanium dioxide is a very well-studied material with wide ranging applications, as discussed in [1].  $TiO_2$  films have been deposited by a wide variety of methods, representatively summarized in [1]. This paper considers  $TiO_2$  as a potentially

useful material for long-wave infrared applications motivated by the relatively low dispersion for  $TiO_2$  in the 8 to 12 µm wavelength range, in comparison to commonlyused SiO<sub>2</sub>. This work presents a potentially advantageous aqueous deposition method that provides smooth conformal coatings. We designate the film considered here as "smooth SPEED"  $TiO_2$  to distinguish it from a highly structured film grown by a different SPEED recipe reported earlier [1], which was designated "Ropy SPEED" according to its observed morphology.

The complex refractive index spectra have been reported out to far infrared wavelengths (125  $\mu$ m) for TiO<sub>2</sub> thin films prepared by vacuum-based physical deposition [2,3]. Our atmospheric-pressure aqueous inhomogeneous-chemical-reaction method has unique potential for low cost, large area, perfectly conformal [1, 4] roll-to-roll manufacturing. Physical and optical properties of such films are reported here with particular emphasis on application to infrared photonics, such as metal-insulator-metal plasmonic resonant absorbers [5] and planer waveguides [6].

#### **METHODS**

Streaming Process for Electrodeless Electrochemical Deposition (SPEED) grows films by heterogeneous reaction, with little wasteful homogeneous reaction, in contrast to chemical bath deposition. Hydrophilic substrates bind hydroxyl ion (OH<sup>-</sup>) nucleation sites with density  $> 10^{12}$  cm<sup>-3</sup>. Substrate heating provides reaction activation energy. Precursors comprised titanium chloride in organic solvents and 20%-by-volume deionized water, which sources OH- and oxygen. Ethanol, isopropanol, and methyl propanol are complexing agents for Ti ions. Films were grown on bare or aluminum-coated silicate glass. Both surfaces are hydrophilic and readily anchor OH<sup>-</sup> ions, which attract positively charged complexes to initiate electrochemical reaction without external field or electrodes. The heterogeneous chemical reactions are

$$[Sub](OH)^{-} + [TiL_n]^{p+(4+)} \rightarrow [Sub]Ti(OH)^{3+} + (L_n)^{p\uparrow}$$
(1)

$$[Sub]Ti(OH)^{3+} + 3OH^{-} \rightarrow [Sub]TiO_2 + 2H_2O\uparrow$$
(2)

where L is an organic ligand, [Sub] the substrate, p the ligand charge, and n the number of ligands coordinated with Ti. Substrate temperature must exceed the reaction activation energy and suffice to eliminate by-products by decomposition and evaporation (up arrows). Site regeneration follows

$$[Sub]TiO_2 + OH^- \rightarrow [Sub]TiO_2(OH)^-.$$
(3)

The freshly attached OH<sup>-</sup> initiates the next cycle. Growth competition leads to nanoparticle morphology.

Asymmetric out-of-plane X-ray Diffraction (PANalytical Empyrean, XRD) used CuK $\alpha$  radiation ( $\lambda$ = 0.154059 nm) at incidence angle 1 deg. XRD data were compared with reference spectra for Anatase, Rutile, and Brookite (International Centre for Diffraction Data, PDF2 Release 2015). Scanning electron microscopy (Zeiss ULTRA-55) imaged surface. Cut and polished cross-sections (Allied Multiprep Polishing System-12 and 10:1 Buffered Oxide Etch) were imaged using scanning ion beam microscopy (Zeiss Cross Beam FIB). Samples for SEM imaging were sputter coated with 2 nm Iridium (Gatan Precision Etching Coating System).

Atomic Force Microscopy (Digital Instruments 5000U, AFM, with Gwyddion software) gave surface roughness map. Transmittance spectra (Cary 500i) determined

band gap. Optical constants were determined by spectroscopic ellipsometry using a J. A. Woollam IR-VASE.

#### RESULTS

Figure 1 (left) presents XRD results, which show that smooth SPEED  $TiO_2$  on Al-coated glass is Anatase phase. Growth of  $TiO_2$  on a conducting substrate was motivated by the LWIR applications presented below.) Aluminum peaks are also observed.

Figure 1 (right) presents SEM images of the top surface of the film. The surface is mostly smooth and featureless, except for small scattered blisters of ~ 1  $\mu$ m lateral dimensions. In the smooth regions, the microstructure has ~20 nm length scale. These films are suitable for long-wave infrared (LWIR) applications, because those wavelengths (~10  $\mu$ m) are much larger than the surface defects. Blisters may form from a combination compressive stress and poor adhesion, noting that the thermal expansion coefficients of aluminum and anatase are significantly different, or these surface defects may be non-volatized by-products.



Figure 1. (left) XRD of TiO<sub>2</sub> film and reference. (right) SEM images of smooth-SPEED TiO<sub>2</sub> film surface.

Figure 2 (left) presents a cross section showing substrate, Al & TiO<sub>2</sub> layers, and 200 nm scale bar. The TiO<sub>2</sub> film is ~ 130 nm thick. Figure 2 (right) shows surface morphology comprising sparse and isolated narrow peaks of ~200 nm amplitude and ~ 1  $\mu$ m wide base on a smooth background of ~10 nm roughness, which is consistent with the SEM image in Figure 1.



Figure 2. (left) SEM cross-sectional images of Smooth SPEED TiO2 film. (right) AFM image of Smooth SPEED TiO2.

Figure 3 (left) presents a UV-visible transmission spectrum for a film on glass, using bare glass reference. The refractive index of Anatase at 2.5 eV is ~3 for unpolarized light and random crystallite orientation [7], giving 25% Fresnel reflectance at 500 nm wavelength. The 80% peak transmittance somewhat exceeds expectations, which may be a due to interference within the film, since the optical thickness n\*d ~ 390 nm exceeds half the 500 nm wavelength. This is supported by the observation of one full Fabry-Perot oscillation for a 210 nm thick evaporated TiO<sub>2</sub> film. For the thinner smooth-SPEED film, Fabry-Perot fringes should be separated by 1.5 eV, and a weak bump appears 1.2 eV below the 3 eV peak. The smaller than expected period may be due to truncation by fundamental absorption above 3 eV.



Figure 3. (left) Transmission spectrum (right) Determination of  $E_g$ .

Figure 3 (right) presents determination of the indirect gap of anatase [7]  $E_g$ . Absorption coefficient  $\alpha$  was calculated according to Beer's law  $\ln(1/T)/d$  from film thickness *d* and the transmittance spectrum *T*, which was normalized to remove Fresnel reflectance. The absorption edge depends on absorption and emission of phonons with energy  $E_{p_2}$  according to [8]

$$\alpha_{\rm e} + \alpha_{\rm a} = C_{\rm e}(T) \left(h\nu - E_g - E_p\right)^2 + C_{\rm a}(T) \left(h\nu - E_g + E_p\right)^2 \quad , \tag{4}$$

where the coefficients depend on temperature-dependent phonon statistics. A plot of the square root of equation 4 gives two straight segments, whose extensions to  $\alpha = 0$  define  $E_g \pm E_p$ . The determined values are given in the Figure. If  $\sqrt{\alpha}h\nu$  is plotted instead [9], somewhat larger values  $E_g = 3.21$  eV and  $E_p = 0.21$  eV are found.  $E_g$  is usually reported

to be in the range 3.08-3.35 eV [7, 10-14]. The  $E_p$  value is about twice the onset of optical phonon absorption, shown below.

At high energies, the absorption coefficient goes as  $\sqrt{(hv - E_g')}$ , where  $E_g'$  is the direct gap. We find  $E_g'$  to be about 4.2 eV by extrapolating a plot of  $\alpha^2$  (or  $(\alpha hv)^2$ ) to zero absorption. This value is higher than the expected value 3.8 eV [7]. However, the spectral region that follows the square root dependence is narrow with very low transmittance, so that our value of  $E_g'$  has high uncertainty.

Figure 4 (left) compares infrared spectra of refractive index n and extinction coefficient  $\kappa$  for smooth SPEED TiO<sub>2</sub> to those for evaporated SiO<sub>2</sub>. The TiO<sub>2</sub> spectra compare well with previous reports for differently prepared films [2,3]. SiO<sub>2</sub> has a strong and sharp extinction peak near 9 µm wavelength, but the TiO<sub>2</sub> extinction is comparatively weak out to at least 11 µm wavelength. Thus, TiO<sub>2</sub> would be better than SiO<sub>2</sub> as an index contrast material for LWIR photonic planer waveguides due to its comparatively low loss [6]. Moreover, TiO<sub>2</sub> lacks the derivative-like dispersion feature observed for n of SiO<sub>2</sub>, which allows TiO<sub>2</sub>-based metal-insulator-metal (MIM) absorbers to have single resonances for spectral sensing applications as opposed to the complicated spectra observed for SiO<sub>2</sub>-based devices [15]. Figure 4 (right) demonstrates the difference for SiO<sub>2</sub>- and SPEED-TiO<sub>2</sub>-based MIM structures with (top-square, dielectric thicknesses) = (3.1 µm, 1.3 µm) and (1.31 µm, 0.104 µm), respectively.

The onset of extinction beyond 12  $\mu$ m wavelength suggests optical phonon energies less than 0.1 eV. This is agreement with optical phonon frequencies determined from far-infrared reflectivity [16]. Thus, the determined E<sub>p</sub> value (Figure 3) is surprisingly large, unless the absorption is a two phonon process.



Figure 4. (left) Infrared complex index spectra for smooth SPEED TiO<sub>2</sub> and evaporated SiO<sub>2</sub> films. The shaded region indicates LWIR wavelengths. (right) Reflectance spectra of SiO<sub>2</sub> and SPEED TiO<sub>2</sub>-based MIM devices.

#### DISCUSSEION AND CONCLUSIONS

Alternative aqueous spray methods of  $TiO_2$  film deposition include liquid tetraisopropyl titanate (TPT) and water separately picked up by inert carrier gas, mixed, then sprayed on heated substrate (100-400 C) [17]. Deposition rate was only 0.1-1 nm/sec. Only optical characterization without microstructural information were

presented. In contrast, SPEED TiO<sub>2</sub> films grew at 100-1000 higher rates by directly spraying an aqueous solution containing titanium ions and ligands onto substrates at 300-400 C.

Spray pyrolysis [18] of organo-metallic compounds to grow TiO<sub>2</sub> requires sufficient heat to break (lyse) precursor molecules, whereas SPEED proceeds via heterogeneous chemical reaction, heat being used mainly to evaporate reaction byproducts.

Spray coating of TiO<sub>2</sub> film can be done at ambient conditions using a suspension of precrystallized nano-particles [19]. Resulting smooth films up to 30 nm thick suffice for electron-transport in solar cells, but this is too thin for infrared applications. Such films are less adhesive than SPEED films, which are chemicallybound to the surface, in our experience.

In conclusion, optical and morphological properties of Smooth Anatase TiO<sub>2</sub> film grown by aqueous spray deposition were presented. Defects are much smaller than long-wave infrared wavelengths and should not contribute significantly to electrodynamic losses. Low dispersion in the long-wave infrared is attractive for resonant absorbers in applications that require low loss and spectral selectivity.

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